

Viscorotational shear instability in astrophysical dense granular flows

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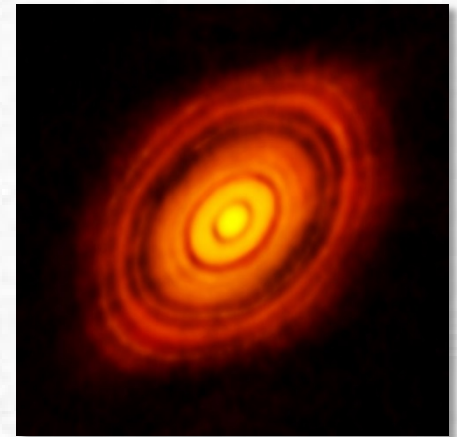
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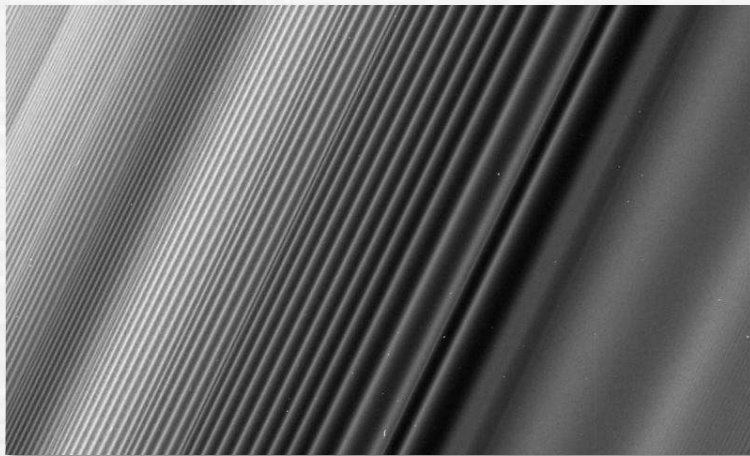
Granular Flows in Astrophysics

- **Planetary Rings**
- **Exo-Planetary Rings**
- **Debris Discs**
- **Protoplanetary Discs**



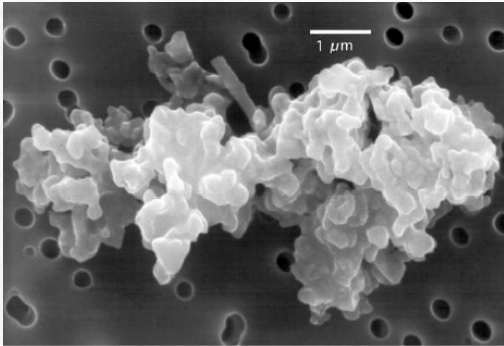
Differentially rotating “Keplerian flows”:

Particles rotate on their gravitational orbits $\Omega(r) \propto r^{-3/2}$



Granular pressure comparable with residual gravitational pressure

Planet Formation



Sizes of solid particles

$1\ \mu\text{m} \rightarrow 10\ \text{km}$

dust \rightarrow planetesimal



Particle collisions: coagulation, scattering, fragmentation

Core accretion theory: ~ 1 meter barrier

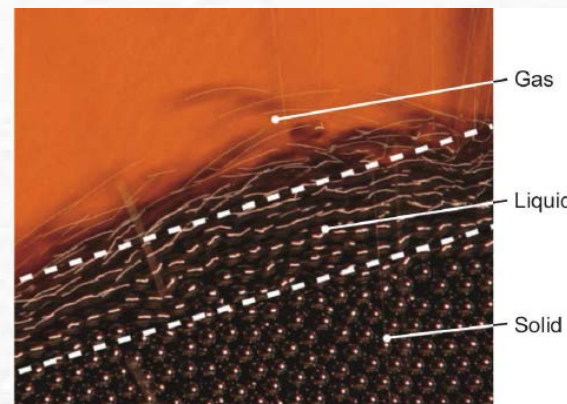
a) Fluid model: nonlinear dynamics of anticycl. vortices;

b) Fluid-Particle model: “streaming instability”;

c) Particle-Particle collisions? granular fluid;

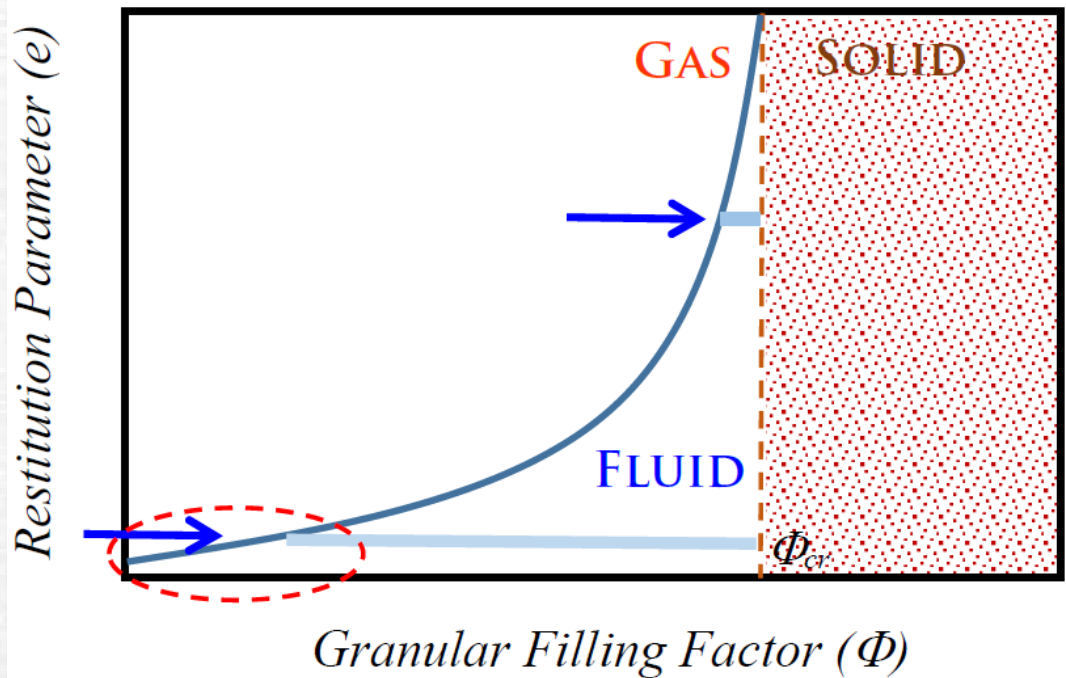
Granular Liquids

Granular flow behavior
Strongly depends on the
Particle restitution coefficient.



Granular liquid
of inelastic particles
can be observed for
a wide range of
particle filling factors

*Granular Rheology
of disc flows*



Viscous Stability

Turbulent (anomalous) viscosity α -model: $\nu_{\text{eff}} = \alpha hc$

Shakura Syunyaev 1973

$$\nu \propto \sigma^\beta$$

σ – surface density (opacity)

a) Viscous Instability: $\beta < -1$ (incompressible)

Lightman, Eardley 1974, Shakura Syunyaev 1976

Viscous stress is proportional to radiation pressure?

b) Viscous Overstability: $\beta > 2$ (compressible)

Kato 1978, Blumenthal et al. 1984

Increase of viscous stress in compressed d-spiral wave;

Dense Granular Flow Model

$$\rho \left\{ \frac{\partial}{\partial t} + V_k \frac{\partial}{\partial x_k} \right\} V_i = -\frac{\partial P}{\partial x_i} + \rho \frac{\partial \Phi_0}{\partial x_i} + \frac{\partial \tau_{ik}}{\partial x_k}$$

- Incompressible flow;
- Constant Gravitational Potential;
- Granular Viscosity;
- Tensor formulation; (*Jop, Forterre & Pouliquen 2006*)

$$\frac{\partial V_k}{\partial x_k} = 0$$

$$\tau_{ij} = \eta \dot{\gamma}_{ik}$$

$$\xi \equiv \sqrt{\dot{\gamma}_{ik} \dot{\gamma}_{ik} / 2}$$

- Local Constitutive Law

$$\eta = \eta(P, \xi)$$

Dense Granular Flow Model

Equations of motion in cylindrical co-ordinates

$$\begin{aligned} \frac{\partial v_\phi}{\partial t} + \left(v_r \frac{\partial}{\partial r} + \frac{v_\phi}{r} \frac{\partial}{\partial \phi} + v_z \frac{\partial}{\partial z} \right) v_\phi + \frac{v_r v_\phi}{r} = & -\frac{1}{\rho r} \frac{\partial P}{\partial \phi} + \frac{1}{r} \frac{\partial \Phi}{\partial \phi} + \\ & + \frac{\eta}{\rho} \left(\Delta v_\phi - \frac{v_\phi}{r^2} + \frac{2}{r^2} \frac{\partial v_r}{\partial \phi} \right) + \frac{1}{\rho} \frac{\partial \eta}{\partial r} \left(\frac{1}{r} \frac{\partial v_r}{\partial \phi} + \frac{\partial v_\phi}{\partial r} - \frac{v_\phi}{r} \right) + \\ & + \frac{2}{r} \frac{\partial \eta}{\partial \phi} \left(\frac{1}{r} \frac{\partial v_\phi}{\partial \phi} + \frac{v_r}{r} \right) + \frac{1}{\rho} \frac{\partial \eta}{\partial z} \left(\frac{\partial v_\phi}{\partial z} + \frac{1}{r} \frac{\partial v_z}{\partial \phi} \right) \end{aligned}$$

Linear Perturbations

Viscosity fluctuations:

$$\frac{\eta'}{\rho} = G_P \frac{P'}{\rho} + G_S \left(\frac{\partial V'_x}{\partial y} + \frac{\partial V'_y}{\partial x} \right)$$

Pressure rheology parameter:

$$G_P = \left(\frac{\partial \eta}{\partial P} \right)_{\xi}$$

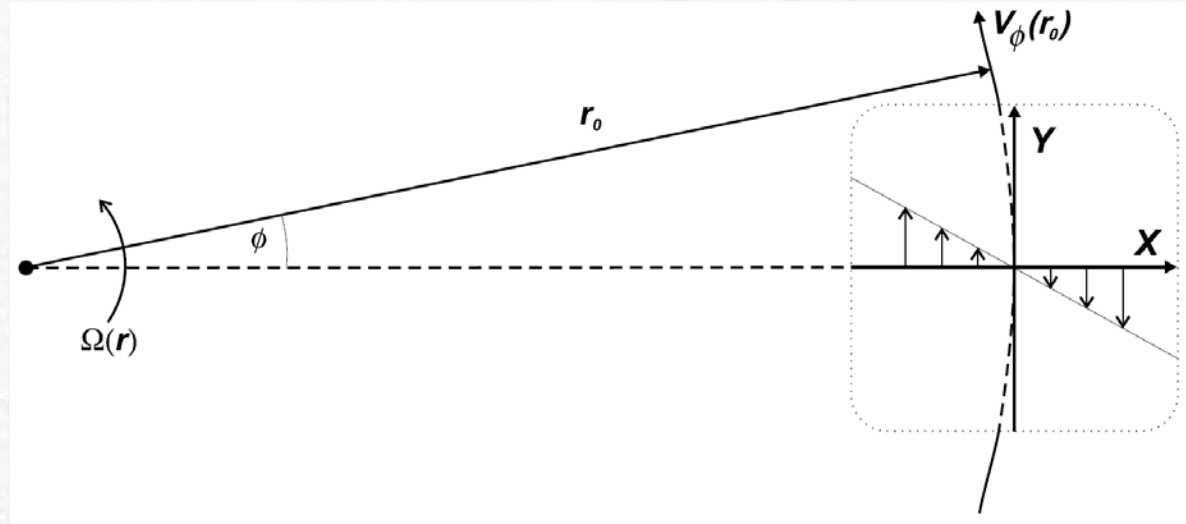
Shear rheology parameter:

$$G_S = \frac{1}{\rho} \left(\frac{\partial \eta}{\partial \xi} \right)_P$$

Linear Perturbation Analysis

Co-rotating
cylindrical frame:

Shearing sheet
(Kelvin modes)



Differential rotation \rightarrow Planar shear flow

Shearing sheet modes

$$\psi'(t) \propto \exp(i\mathbf{k}(t)\mathbf{r} - i\omega t)$$

$$k_x(t) = k_0 - 2Ak_x t$$

$$\omega(t) = \omega(\mathbf{k}(t))$$

Linear Stability Analysis

Solutions of the dispersion equation:

$$\omega(\mathbf{k}) = \pm \left(\kappa^2 - W^2 \right)^{1/2} + i(W - \nu k^2)$$

Epicyclic frequency in Newtonian fluids:

$$\kappa_0^2 = \frac{2\Omega_K(r)}{r} \frac{d}{dr} \left(r^2 \Omega_K(r) \right)$$

Rheological modification:

$$\kappa^2 = \kappa_0^2 + \Delta K^2(G_P, G_S)$$

and

$$W = \sigma_A + \sigma_P + \sigma_S$$

Linear Stability Analysis

Transient growth:

$$\sigma_A = \frac{Ak_x k_y}{k^2 - 4AG_P k_x k_y}$$

Pressure rheology effects:

$$\sigma_P = 2AG_P \frac{(\Omega k_x^2 + Bk_y^2)}{k^2 - 4AG_P k_x k_y}$$

Shear (strain) effects:

$$\sigma_S = -AG_S \frac{(k_x^2 - k_y^2)^2 + k_{\perp}^2 k_z^2}{k^2 - 4AG_P k_x k_y}$$

Epicyclic modification:

$$\Delta K^2(G_P, G_S) = -4G_S \frac{A^2 k_x k_y k_z^2}{k^2 - 4AG_P k_x k_y}$$

Instability

Linear instabilities:

Case 1: Pressure rheology:

$$G_P < 0, \quad G_S = 0$$

“viscous instability” (cf. $\beta < -1$)

Case 2: Shear rheology:

$$G_P = 0, \quad G_S > 0$$

$$W^2 > \kappa^2$$
$$W + \sqrt{W^2 - \kappa^2} > \nu k^2$$

- Keplerian flow: $\xi(r_0) = -2A$

- Power law constitutive eq.:

$$\left(\frac{\ln \eta}{\ln \xi} \right)_P > 2$$

Visco-rotational Shear Instability

- Vertically uniform perturbations: $k_z = 0$
- Shear rheology: $G_P = 0$, $G_S \neq 0$
- Arbitrary differential rotation: $\Omega(r) \propto r^{-q}$

$$\frac{d}{dt} (\ln(\text{curl} \mathbf{v})_z) = q G_S \Omega \frac{(k_x(t)^2 - k_y^2)^2}{k(t)^2} - \nu k(t)^2$$

*radial
differential
rotation*

*shear
rheology*

*Angular
velocity of
rotation*

- 1)** $q > 0$, $G_S > 0$; ($q_K = 3/2$)
- 2)** $q < 0$, $G_S < 0$;

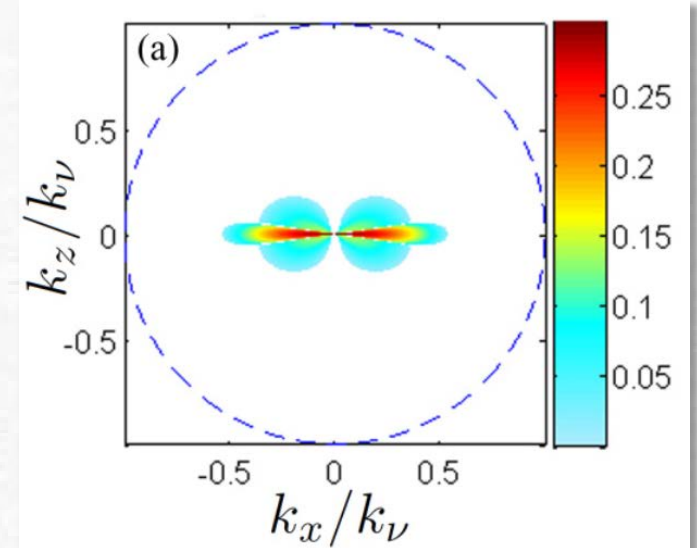
Instability Spectra

1) Pressure rheology instability:

$$G_P \neq 0, \quad G_S = 0$$

Large scale

“Viscous instability”

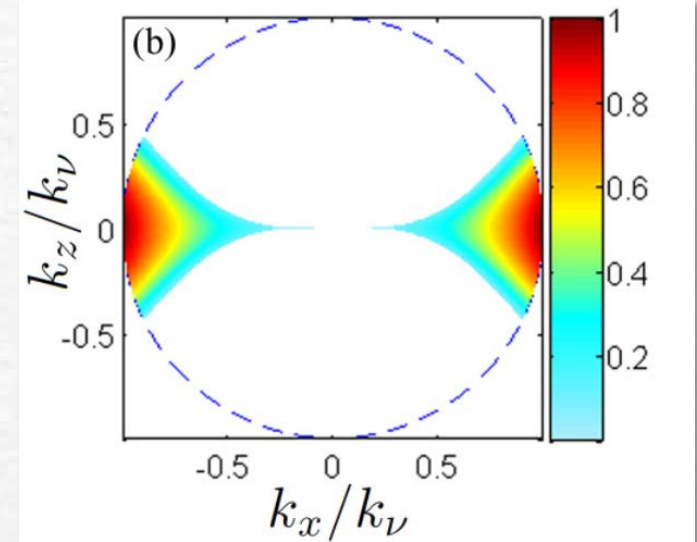


2) Shear rheology instability:

$$G_P = 0, \quad G_S \neq 0$$

Small scale

“Visco-rotational instability”



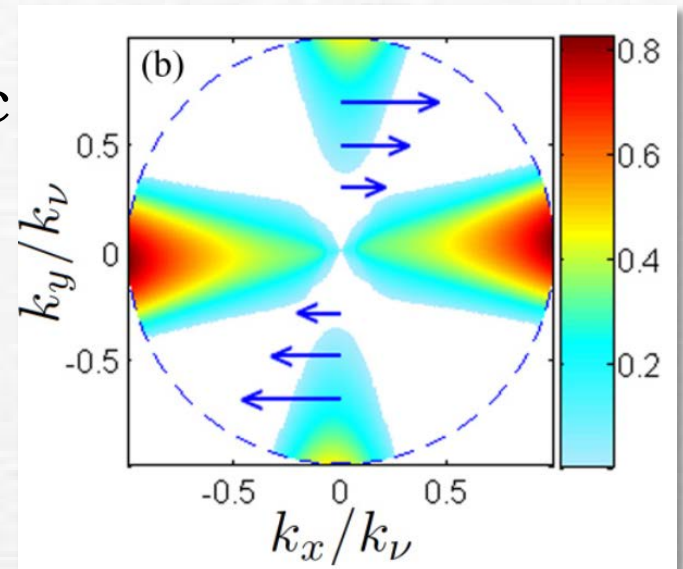
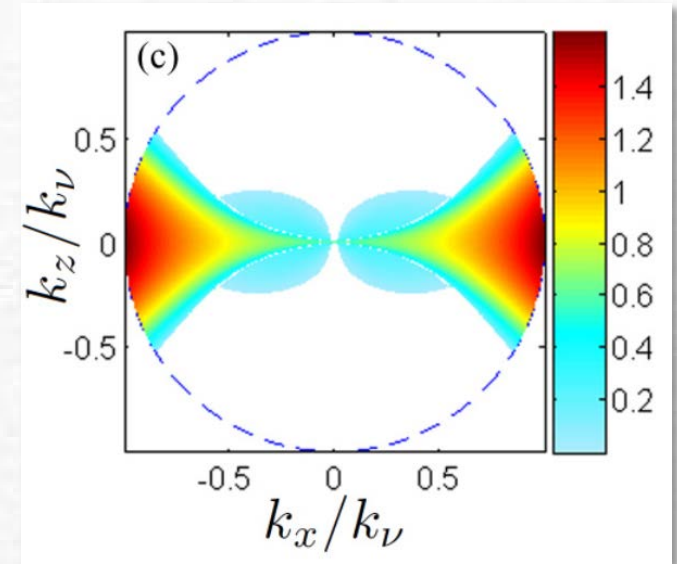
Instability Spectra

3) Instability of disk flow with pressure and shear rheologies:

$$G_P \neq 0, \quad G_S \neq 0$$

4) Instability of non-axisymmetric Perturbations:

$$k_y \neq 0, \quad k_x = k_x(t)$$



Summary

Shear and Pressure rheology effects on linear waves in Keplerian discs.

New type of linear instability:
“visco-rotational shear” instability

$$\left(\frac{\ln \eta}{\ln \xi} \right)_P > 2$$

- ✓ Incompressible, shear rheology,
- ✓ Epicyclic destabilization of perturbations;
- ✓ Exponential growth of the perturbation velocity curl in the shear plane;
- *Constitutive law dependence of particle sizes?*
- *Nonlinear development: numerical simulations?*

Thank You

Equilibrium Disc Flow

$$r\Omega^2 = -\frac{\partial\Phi}{\partial r},$$
$$\left(r\frac{\partial^2\Omega}{\partial r^2} + 3\frac{\partial\Omega}{\partial r}\right)\bar{\eta} + r\frac{\partial\Omega}{\partial r}\frac{\partial\bar{\eta}}{\partial r} = 0,$$

$$\Phi(r,z) = \frac{GM}{(r^2 + z^2)^{1/2}}$$

$$\Omega(r) = \Omega_0\left(\frac{r}{r_0}\right)^{-q}, \quad \Omega_0 = \left(\frac{GM}{r_0^3}\right)^{1/2}.$$

$$\frac{\partial \ln \bar{\eta}}{\partial \ln r} = q - 2.$$